

USE OF CREEPING SPEED MODE IN DISCRETE CONTROL AND MEASURING EQUIPMENT WITH DIGITAL-PROGRAM CONTROL

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Abstract. The creeping speed mode is studied with calculation of the resultant vector of the flux linkage ψ_0 of the double-fed machine. It is shown that use of the double-fed machine as the actuating motor in the creeping speed mode increases the working off accuracy of the motion law. A mathematical model of the double-fed machine to study various modes of the CNC machine including the creeping speed mode is developed. The method to estimate the load parameters of the double-fed machine is presented. Practical implication of the study is confirmed by the results obtained in the evaluation of the influence of load parameters on the dynamic indicators of actuating electromotor.

1 Introduction

Currently, any measuring equipment incorporates the systems of the fixed current parameters, comparison with a reference element and compensation of the measuring error [1]. In these systems, modern CNC machines are used; many engines with relatively small power are involved in their work, the basic modes of operation of these engines are “unnatural” [2]. These modes include frequent reversal, braking by various ways (dynamic, double-current, capacitor-dynamic, etc.), getting the creeping speed according to the schemes for the joint feeding, step and torque modes, etc. [3–9].

The creeping speed mode is described in detail below. The feature of the creeping speed mode is an operation at the speed of the tool movement from 2 to 10 mm/min [3]. This mode is widely used in the CNC machines (for processing and plasma cutting), aboard mechanisms of antenna rotation, detection systems for mobile and ground equipment, fixation systems of outdoor surveillance, textile industry (warp sizing machine), and discrete measurement systems with CNC (two-dimensional electromagnetic displacement sensors) [10–12]. In practice, this type of the tool movement is often implemented due to the presence of linear asynchronous engines (LAE), actuating motor and ball-and-screw unit or electromagnetic frame.

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2 Research method

The method to develop the creeping speed in the double-fed machine (DFM) at the different inclusions of magnetic fields is presented. The analysis of influence of the load parameters on the dynamic indicators of actuating electromotor is carried out.

The double-fed machine is referred to as the controlled (full controlled) machine of alternating current; voltages are supplied to its stator and rotor clips. These voltages are composed of two components: voltage of compensation and voltage (law) of control. This type of electrical machines is known to be universal; it has high power characteristics. This electrical machine is implemented as a micromachine; it is a decisive factor for measurement systems [13, 14].

The creeping speed mode (Figure 1) is implemented in the feeding of stator and rotor windings of the DFM by sinusoidal voltage of different frequencies. In the first case, the difference between frequencies is 1 Hz; in the second case, it is 40 Hz at the initial phase of the supply voltage $\varphi=\pi/2$.

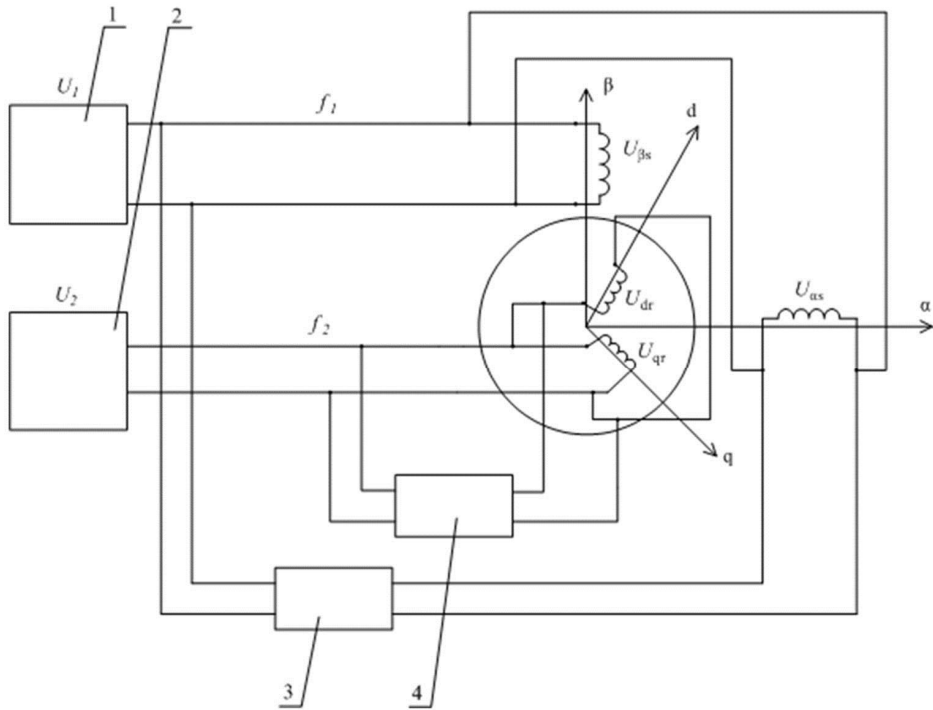


Figure 1. Feeding scheme of the double-fed machine in realization of the creeping speed mode: 1, 2 are the sources; 3, 4 are the phase shifters.

The following equations are true for operation of the DFM in the creeping speed mode: $\omega_1=\omega_2=\omega$, and $\omega_3=\omega_4=\omega'$ if $\alpha=\beta=\gamma=0$, $\varphi=\pi/2$.

Equations (1) for the flux linkages of stator and rotor windings in the transformed coordinate system with due regard to the accepted assumptions and the mode of the machine operation are

$$\begin{aligned}
\psi_{\alpha s} &= \psi_{ms} \sin(\omega t); \\
\psi_{\beta r} &= \psi_{ms} \sin(\omega t); \\
\psi_{\alpha r} &= \psi_{mr} \sin(\omega t) \cos \chi + \psi_{mr} \sin\left(\omega t + \frac{\pi}{2}\right) \sin \chi; \\
\psi_{\beta r} &= -\psi_{mr} \sin(\omega t) \sin \chi + \psi_{mr} \sin\left(\omega t + \frac{\pi}{2}\right) \cos \chi.
\end{aligned} \tag{1}$$

Equations (2) for the squares of the real and imaginary parts of the generalized flux linkage vector are

$$\begin{aligned}
\psi_{\alpha}^2 &= \left(\psi_{ms} \sin(\omega t) + \left(\psi_{mr} \sin(\omega t) \cos \chi + \psi_{mr} \cos(\omega t) \sin \chi \right) \right)^2 = \\
&= \left(\psi_{ms} (\sin(\omega t) + \mu (\sin(\omega t) \cos \chi + \cos(\omega t) \sin \chi)) \right)^2 = (\psi_{ms} (\sin(\omega t) + \mu \sin(\omega t + \chi)))^2, \\
\psi_{\beta}^2 &= \left(\psi_{ms} \sin(\omega t) - \left(\psi_{mr} \sin(\omega t) \sin \chi + \psi_{mr} \cos(\omega t) \cos \chi \right) \right)^2 = \\
&= \left(\psi_{ms} (\sin(\omega t) - \mu (\sin(\omega t) \sin \chi + \cos(\omega t) \cos \chi)) \right)^2 = (\psi_{ms} (\sin(\omega t) - \mu \cos(\omega t - \chi)))^2.
\end{aligned} \tag{2}$$

Equations (3) for the amplitude value of the flux linkage vector in the creeping speed mode are

$$\psi = \psi_{ms} \sqrt{(\sin(\omega t) + \mu \sin(\omega t + \chi))^2 + (\sin(\omega t) - \mu \cos(\omega t - \chi))^2}. \tag{3}$$

Equations (2) for the movement of the spatial vector of flux linkage in the air gap for the creeping speed mode (3) are

$$\begin{aligned}
\chi_0 &= \arctg \left(\frac{\psi_{ms} \sin(\omega t) - \psi_{mr} \sin(\omega t) \sin \chi + \psi_{mr} \cos(\omega t) \cos \chi}{\psi_{ms} \sin(\omega t) + \psi_{mr} \sin(\omega t) \cos \chi + \psi_{mr} \cos(\omega t) \sin \chi} \right) = \\
&= \arctg \left(\frac{\sin(\omega t) - \mu (\sin(\omega t) \sin \chi + \cos(\omega t) \cos \chi)}{\sin(\omega t) + \mu (\sin(\omega t) \cos \chi + \cos(\omega t) \sin \chi)} \right) = \\
&= \arctg \left(\frac{\sin(\omega t) - \mu \sin(\omega t + \chi)}{\sin(\omega t) + \mu \cos(\omega t - \chi)} \right).
\end{aligned} \tag{4}$$

Figure 2 shows hodographs of the resultant flux linkage vector at the feeding of stator and rotor windings from the current and voltage sources. Two rotating electromagnetic fields arise during the DFM operation in the creeping speed mode for hodograph of the resultant vector of flux linkage ψ_0 : the first field is generated by the frequency $f_1=f_2=50$ Hz and the second one is generated by $f_3=f_4=51$ Hz from the primary and secondary elements respectively.

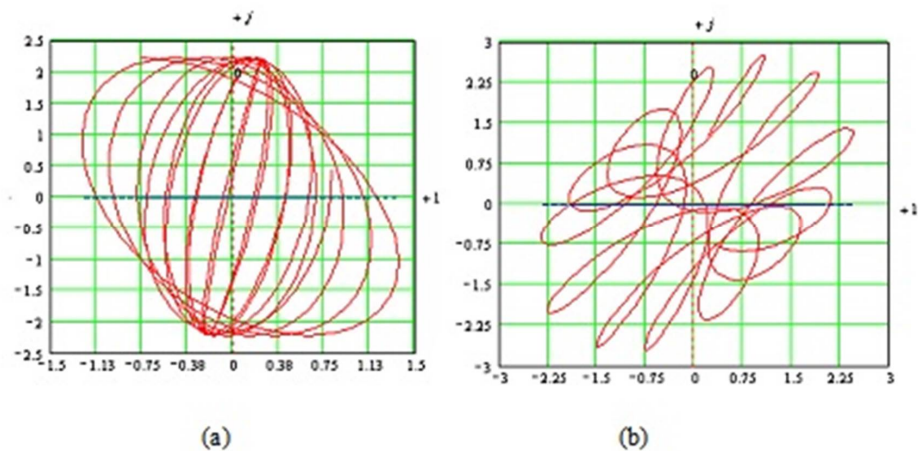


Figure 2. Hodographs of vectors ψ_0 at the feeding of stator and rotor windings from the current source (a) and from the voltage source (b) at $\omega_1=\omega_2=1$; $\omega_3=\omega_4=1.02$.

Consider two variants of interaction of the magnetic fields in the DFM in the creeping speed mode:

1. It is necessary to set the following signals of the control system at the consonant interaction of electromagnetic fields: $\gamma_1=\gamma_2=1$; $\gamma_3=\gamma_4=1.225$; $\omega_1=\omega_2=\omega$; $\omega_3=\omega_4=\omega'$; $\alpha=\beta=\gamma=0$, $\varphi=\pi/2$. The voltages on the clips of stator and rotor windings in transformed axes α , β , 0 are in the following form:

$$\begin{aligned} U_{\alpha s} &= U_{ms} \sin(\omega t); \\ U_{\beta s} &= U_{ms} \sin(\omega t); \\ U_{\alpha r} &= U_{mr} \sin(\omega' t) \cos \chi + U_{mr} \sin(\omega' t + \phi) \sin \chi; \\ U_{\beta r} &= -U_{mr} \sin(\omega' t) \sin \chi + U_{mr} \sin(\omega' t + \phi) \cos \chi. \end{aligned} \tag{5}$$

2. In the case of the opposite interaction of electromagnetic fields the angular frequency of rotation and the initial phases on the axes are set similarly to the above case $\gamma_1=\gamma_2=1$; $\gamma_3= 1.225$; $\gamma_4= -1.225$. Equations for the voltages on the clips of the stator and rotor windings are similar to the system (5) with the only difference: the sign "+" before the second summand is replaced for the sign "-" in the second part of the two equations.

A mathematical model of DFM is given relative to the currents of stator and rotor in the axes α , β , 0 in the iterative form at zero initial conditions. It is presented in the system (6). DFM was based on the asynchronous engine with phase rotor 4AK160S8UZ with protection degree of IP44, and its asynchronous rotational speed is $n=750$ rpm with power $P_2=5.5$ [kW].

This model allows investigating transients in the DFM in various modes of its operation. Table 1 shows the parameters of actuating motor 4AK160S8UZ.

Table 1. Parameters of actuating motor 4AK160S8UZ.

Energy indicators			Mechanical characteristic			Equivalent circuit parameters				
$\cos \varphi$	p	η	m_m	S_{nom}	S_{crit}	X_μ	R'_1	X'_1	R'_2	X'_2
[pu]	[-]	[%]	[pu]	[%]	[%]	[pu]	[pu]	[pu]	[pu]	[pu]
0.7	4	80	0.5	6.4	9	1.6	0.06	0.11	0.094	0.175

Gabriel Kron's equations are the basis for the DFM mathematical model. Gabriel Kron's equations are the system of differential equations of voltages and motions with due regard to the accepted assumptions and describe the energy transformations in the electric machine. It is also necessary to calculate the base values and coefficients of the model, which are shown in Tables 2 and 3.

Table 2. Estimated base values of parameters of actuating motor.

Main base values				Auxiliary base values						
U_b	I_b	P_b	ω_b	t_b	Z_b	L_b	$M_{\text{mech. b}}$	$L_{\text{mech. b}}$	$R_{\text{mec h. b}}$	$C_{\text{mech. b}}$
[V]	[A]	[W]	[rad/s]	[s]	[Om]	[H]	[W·s/rad]	[W·s ³ /rad ³]	[W·s ² /rad ²]	[W·s/rad]
311.1	18.2	8.5	314.3	0.0032	17.07	0.054	108.297	0.00439	1.38	433.19

Table 3. Coefficients of mathematical model [pu].

σ	α_s	α'_s	α_r	α'_r	k_s	k_r
0.158	3.702	0.222	3.756	0.213	0.935	0.901

Mathematical model of DFM is of the form [pu]:

$$\begin{pmatrix} t_{j+1} \\ i_{s\alpha_{j+1}} \\ i_{s\beta_{j+1}} \\ i_{r\alpha_{j+1}} \\ i_{r\beta_{j+1}} \\ \omega_{j+1} \\ \chi_{j+1} \end{pmatrix} = \begin{pmatrix} t_j + dt \\ i_{s\alpha_j} + dt \cdot \left(\alpha_s U_1 \sin(\omega_1 t_j + \alpha) - \alpha_s K_R (U_3 \sin(\omega_1 t_j + \gamma) \cdot \cos \chi_j \dots + U_4 \sin(\omega_2 t_j + \varphi) \cdot \sin \chi \right. \\ \left. - i_{s\beta_j} \omega_j \frac{K_S \cdot K_R}{\sigma} - i_{s\alpha_j} \dot{\alpha}_s - i_{r\beta_j} \omega_j \frac{K_S}{\sigma} + i_{r\alpha_j} \dot{\alpha}_R K_S \right) \\ i_{s\beta_j} + dt \cdot \left(\alpha_s U_2 \sin(\omega_1 t_j + \beta) - \alpha_s K_R (-U_3 \sin(\omega_1 t_j + \gamma) \cdot \sin \chi_j \dots + U_4 \sin(\omega_2 t_j + \varphi) \cdot \text{cc} \right. \\ \left. - i_{s\alpha_j} \omega_j \frac{K_S K_R}{\sigma} - i_{s\beta_j} \dot{\alpha}_s + i_{r\alpha_j} \omega_j \frac{K_S}{\sigma} + i_{r\beta_j} \dot{\alpha}_R K_S \right) \\ i_{r\alpha_j} + dt \cdot \left(\alpha_R (U_3 \sin(\omega_1 t_j + \gamma) \cdot \cos \chi_j \dots + U_4 \sin(\omega_2 t_j + \varphi) \cdot \sin \chi_j) \dots - K_R \alpha_s U_1 \sin(\omega_1 t_j + \right. \\ \left. + \alpha_s K_R i_{s\alpha_j} - i_{s\beta_j} \omega_j \frac{K_R}{\sigma} + i_{r\beta_j} \omega_j \frac{1}{\sigma} - i_{r\alpha_j} \dot{\alpha}_R \right) \\ i_{r\beta_j} + dt \cdot \left(\alpha_R \cdot (-U_3 \sin(\omega_1 t_j + \gamma) \cdot \sin \chi_j \dots + U_4 \sin(\omega_2 t_j + \varphi) \cdot \cos \chi_j) \dots - K_R \alpha_s U_2 \sin(\omega_2 t_j + \right. \\ \left. - \frac{K_R}{\sigma} \cdot \omega_j \cdot i_{s\alpha_j} - \alpha_s \cdot K_R \cdot i_{s\beta_j} - i_{r\alpha_j} \cdot \omega_j \cdot \frac{1}{\sigma} - i_{r\beta_j} \cdot \dot{\alpha}_R \right) \\ \omega_j + dt \cdot \left(\frac{1}{J} (-L_m (i_{s\beta_j} i_{r\alpha_j} - i_{s\alpha_j} i_{r\beta_j}) - R_g \omega_j - C_m \chi_j - M_c) \right) \\ \chi_j + dt \cdot \omega_j \end{pmatrix} \quad (6)$$

where R_g is a coefficient of damping torque of load; C_m is a coefficient of position torque of load; J is a total torque of inertia of engine; M_c is a torque of static load on the shaft of engine; χ is a law of motion of the movable element.

3 Results and discussion

Figure 3 shows results of mathematical modelling of the DFM (6). As an example, the dependencies of the DFM electromagnetic torque in the starting and creeping speed modes are presented.

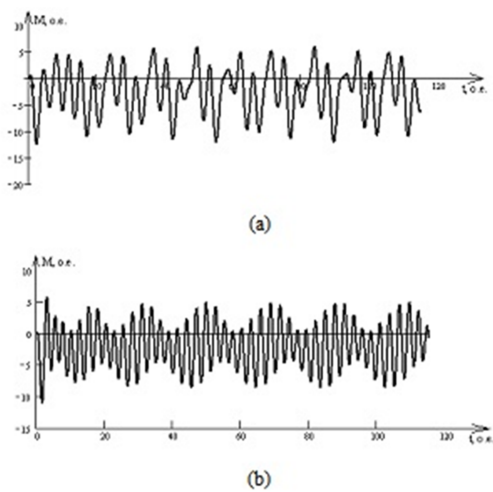


Figure 3. Dependence of electromagnetic moment on the time at the frequency difference of 1 Hz (a) and of 40 Hz (b).

The dependencies were obtained by the modelling; they are adequate and the model is sustainable. The disadvantage is inability to investigate the model in case of simultaneous changes in the load parameters.

The influence of load parameters on the dynamic indicators of actuating motor is important as the additional disturbances and measurement errors lead to incorrect measurement results. The load torque of inertia J_n (from 80 to 2500 pu) is selected as a changeable load parameter.

Figure 4 shows the results of investigation of dynamic indicators: time of transition process τ_{tp} and shock values of electromagnetic torque M_{sh} on load torque of inertia J_n .

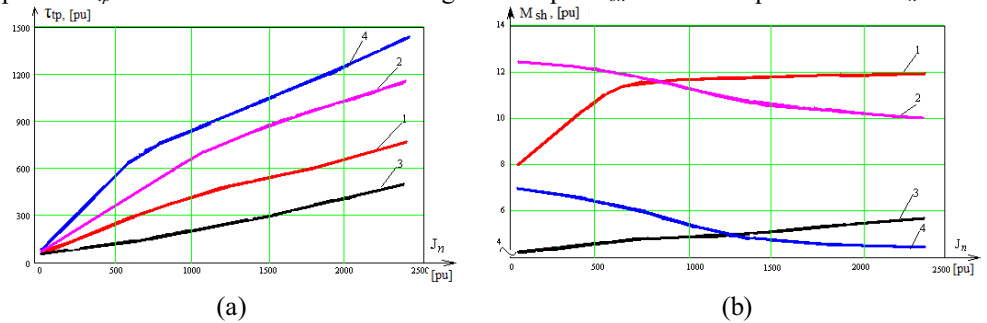


Figure 4. Dependence of time of transition process (a) and shock values of electromagnetic torque (b) on inertia load: consonant interaction of electromagnetic fields at the frequency difference of 1 Hz (curve 1) and of 40 Hz (curve 3); opposite interaction of electromagnetic fields of 1 Hz (curve 2) and of 40 Hz (curve 4).

In the creeping speed mode, the change in time of the transition process on load torque of inertia $\tau_{tp}=f(J_n)$ is observed to increase; it is typical of special modes of electromechanical energy converters based on the DFM.

In the case of the consonant interaction of electromagnetic fields, the shock values of electromagnetic torque M_{sh} decrease with the increase in the inertia torque of load J_n (Figure 4, b , curves 1 and 3), and these values increase in the case of the opposite interaction of electromagnetic fields (Figure 4, b , curves 2 and 4).

4 Conclusion

1. The change of the load torque of inertia J_n leads to the increase in time of transition process τ_{tp} and decrease in the shock values of electromagnetic torque M_{sh} in the case of the consonant interaction of electromagnetic fields. With the same load parameters, the DFM has the maximal time of transition process at the consonant interaction of electromagnetic fields and at the frequency difference of 40 Hz.
2. The practical application of the DFM in discrete-measuring systems of the CNC machines is possible with some limitations: the frequency difference must not exceed 10 Hz.

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